

Half-wave plate systematics: impact on cosmic birefringence and component separation

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searching for *B*-modes from inflation



B-modes can probe inflation.

Unprecedented sensitivity requirements!

the HWP: reducing systematics



A rotating half-wave plate (HWP) as first optical element:

- modulates the signal to $4f_{HWP}$, allowing to "escape" 1/f noise;
- makes possible for a single detector to measure polarization, reducing pair-differencing systematics.

Mueller calculus: radiation described as S = (I, Q, U, V) and HWP effects parametrized by \mathcal{M}_{HWP} , so that $S' = \mathcal{M}_{HWP}S$.



how does this affect the observed maps?

steps we took in that direction

- produce output maps from beamconv-based TOD simulations;
- derive analytic formulae to interpret the output;
- discuss how the non-idealities affect cosmic birefringence.

PREPARED FOR SUBMISSION TO JCAP Impact of half-wave plate systematics on the measurement of cosmic birefringence from CMB Marta Monelli," Eiichiro Komatsu," Alexandre Adler," Matteo polarization Marta Monein, Eucnro Komatsu, Alexanore Adrea Billi 443 Paolo Campeti, 49 Nadia Dachlythra, Adriaan பார், எல்ல கோற்சபு, அவர்க் பகரைபாகு, Adriaan Duivenvoorden,¹, Jon Gudmundsson,^c and Martin Reinecke.^a

Duivenvoorden et al. (2021) MNRAS 502; Monelli et al. arXiv:2211.05685

simulations

Working assumptions: no noise, single freq., CMB-only, simple beams.

- *I*, *Q* and *U* input maps (n_{side} = 512) from best-fit 2018 Planck power spectra;
- 1 year of LiteBIRD-like scanning strategy (mimicking pyScan).
- Instrument specifics: 160 detectors from the 140 GHz channel of LiteBIRD's MFT.
- Non-ideal HWP: Mueller matrix elements from Giardiello et al. (2022) A&A 658.

specs.	values
f _{samp}	19 Hz
HWP rpm	39
FWHM	30.8 arcmin
offset quats.	[]
offset quats.	30.8 arcmin []

ideal and non-ideal TODs, both processed with ideal map-maker.

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(beam transfer function not deconvolved)

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- ▶ *BB* much larger (*EE* shape)

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- TT slightly affected
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- TT slightly affected
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- TE slightly affected
- ► EB non-zero!

ideal and non-ideal TODs, both processed with ideal map-maker.



- TT slightly affected
- EE lost power
- BB much larger (EE shape)
- TE slightly affected
- ► EB non-zero!
- ► TB non-zero!

how can we understand this?

(minimal) TOD: signal detected by 4 detectors. map-maker: bin-averaging assuming ideal HWP. estimated output maps: linear combination of $\{I, Q, U\}_{in}$.

for good coverage and rapidly spinning HWP:

$$\widehat{\mathsf{S}}\simeq egin{pmatrix} m_{ii}\,I_{
m in}\ [(m_{qq}-m_{uu})Q_{
m in}+(m_{qu}+m_{uq})U_{
m in}]/2\ [-(m_{qu}+m_{uq})Q_{
m in}+(m_{qq}-m_{uu})U_{
m in}]/2 \end{pmatrix}.$$

Expanding \widehat{S} in spherical harmonics:

$$\begin{split} \widehat{C}_{\ell}^{TT} &\simeq m_{ii}^{2} C_{\ell,\text{in}}^{TT}, \\ \widehat{C}_{\ell}^{EE} &\simeq \frac{(m_{qq} - m_{uu})^{2}}{4} C_{\ell,\text{in}}^{EE} + \frac{(m_{qu} + m_{uq})^{2}}{4} C_{\ell,\text{in}}^{BB} + \frac{(m_{qq} - m_{uu})(m_{qu} + m_{uq})}{4} C_{\ell,\text{in}}^{EB}, \\ \widehat{C}_{\ell}^{BB} &\simeq \frac{(m_{qq} - m_{uu})^{2}}{4} C_{\ell,\text{in}}^{BB} + \frac{(m_{qu} + m_{uq})^{2}}{4} C_{\ell,\text{in}}^{EE} - \frac{(m_{qq} - m_{uu})(m_{qu} + m_{uq})}{4} C_{\ell,\text{in}}^{EB}, \\ \widehat{C}_{\ell}^{TE} &\simeq \frac{m_{ii}(m_{qq} - m_{uu})}{2} C_{\ell,\text{in}}^{TE} + \frac{m_{ii}(m_{qu} + m_{uq})^{2}}{2} C_{\ell,\text{in}}^{EB}, \\ \widehat{C}_{\ell}^{EB} &\simeq \frac{(m_{qq} - m_{uu})^{2} - (m_{qu} + m_{uq})^{2}}{4} C_{\ell,\text{in}}^{EB} - \frac{(m_{qq} - m_{uu})(m_{qu} + m_{uq})}{4} (C_{\ell,\text{in}}^{EE} - C_{\ell,\text{in}}^{BB}), \\ \widehat{C}_{\ell}^{TB} &\simeq \frac{m_{ii}(m_{qq} - m_{uu})}{2} C_{\ell,\text{in}}^{TB} - \frac{m_{ii}(m_{qu} + m_{uq})}{2} C_{\ell,\text{in}}^{TE}. \end{split}$$

analytical vs non-ideal output spectra



impact on cosmic birefringence

a side effect: measuring cosmic birefringence

CMB might also carry information about parity-violating new physics: cosmic birefringence.

(time-dependent parity-violating pseudoscalar field)



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$$\begin{aligned} a^E_{\ell m, \text{obs}} &= a^E_{\ell m} \cos 2\beta - a^B_{\ell m} \sin 2\beta, \\ a^B_{\ell m, \text{obs}} &= a^E_{\ell m} \sin 2\beta + a^B_{\ell m} \cos 2\beta. \end{aligned}$$
$$\begin{aligned} C^{EB}_{\ell, \text{obs}} &= (C^{EE}_{\ell} - C^{BB}_{\ell}) \sin 4\beta/2 \\ &+ C^{EB}_{\ell} \cos 4\beta. \end{aligned}$$

a side effect: measuring cosmic birefringence

CMB might also carry information about parity-violating new physics: cosmic birefringence.

(time-dependent parity-violating pseudoscalar field)



$$a_{\ell m,obs}^{E} = a_{\ell m}^{E} \cos 2\beta - a_{\ell m}^{B} \sin 2\beta,$$

$$a_{\ell m,obs}^{B} = a_{\ell m}^{E} \sin 2\beta + a_{\ell m}^{B} \cos 2\beta.$$

$$C_{\ell,obs}^{EB} = (C_{\ell}^{EE} - C_{\ell}^{BB}) \sin 4\beta/2 + C_{\ell}^{EB} \cos 4\beta.$$
From *Planck* data:

$$\beta = 0.35 \pm 0.14^{\circ} \text{at } 68\% \text{ C.L.}$$

Constraint expected to improve.

image credit: Yuto Minami; Minami and Komatsu (2020) Phys. Rev. Lett. 125



Degeneracy with cosmic birefringence and polarization angle miscalibration!

In first approximation, HWP induces an additional miscalibration.

This doesn't mean that the HWP will keep us from measuring β , but it shows how important it is to carefully calibrate \mathcal{M}_{HWP} .

including frequency dependence

how does the map-model change

ν

Without HWP:
$$\begin{pmatrix} l_j \\ Q_j \\ U_j \end{pmatrix} = \sum_{\lambda} g_{\lambda} \begin{pmatrix} l_{\lambda} \\ Q_{\lambda} \\ U_{\lambda} \end{pmatrix} + n,$$

With HWP: $\begin{pmatrix} l_j \\ Q_j \\ U_j \end{pmatrix} = \sum_{\lambda} \begin{pmatrix} g_{\lambda}^{ii} & 0 & 0 \\ 0 & g_{\lambda}^{qq-uu} & g_{\lambda}^{qu+uq} \\ 0 & g_{\lambda}^{qq-uu} & g_{\lambda}^{qq-uu} \end{pmatrix} \begin{pmatrix} l_{\lambda} \\ Q_{\lambda} \\ U_{\lambda} \end{pmatrix} + n,$
where $g_{\lambda} = \frac{\int d\nu \ G(\nu) S_{\lambda}(\nu)}{\int d\nu \ G(\nu)}, \quad g_{\lambda}^{ii} = \frac{\int d\nu \ G(\nu) m_{ii}(\nu) S_{\lambda}(\nu)}{\int d\nu \ G(\nu)},$ and so on

HWP non-idealities contribute to gain, polarization-efficiency and cross-polarization leakage.

effective SEDs



$$\sum_{\lambda} \begin{pmatrix} g_{\lambda}^{ii} & 0 & 0 \\ 0 & g_{\lambda}^{qq-uu} & g_{\lambda}^{qu+uq} \\ 0 & g_{\lambda}^{qu+uq} & g_{\lambda}^{qq-uu} \end{pmatrix} \begin{pmatrix} I_{\lambda} \\ Q_{\lambda} \\ U_{\lambda} \end{pmatrix} + n,$$

- Since all these effects are frequency dependent, they affect each component differently,
 - An imprecise calibration of M_{HWP} can lead to complications in the component separation step.

- we are now provided with a simulation pipeline that can be easily adapted to study more complex problems (adding noise, more realistic beams...);
- the analytical formulae represent an alternative tool to study the same problems more effectively (but approximately);
- obvious application: exploiting the analytical formulae to derive calibration requirements for the HWP Mueller matrix elements, so that they don't prevent us from measuring cosmic birefringence, nor spoil the foreground cleaning procedure.

backup

TOD: signal detected by 4 detectors looking at the same pixel;

Detected signal modeled as $d = (1 \ 0 \ 0) \cdot \mathcal{M}_{det} \mathcal{R}_{\xi - \phi} \mathcal{M}_{HWP} \mathcal{R}_{\phi + \psi} \cdot S$;



Apply a bin-averaging (ideal) map-maker to those 4 measurements.

$$\widehat{S} \simeq \begin{pmatrix} m_{ii} l_{in} \\ [(m_{qq} - m_{uu})Q_{in} + (m_{qu} + m_{uq})U_{in}]/2 \\ [-(m_{qu} + m_{uq})Q_{in} + (m_{qq} - m_{uu})U_{in}]/2 \end{pmatrix}$$

 $\theta_{EB}, \ \theta_{TB} \ \text{and} \ \widehat{\theta}$

analytical expectation: $\hat{\theta} \sim 3.8^{\circ}$. compatible with best fit estimates!

